

# GaN Microwave Electronics

Umesh K. Mishra, *Fellow, IEEE*, Yi-Feng Wu, Bernd P. Keller, Stacia Keller, and Steven P. Denbaars

**Abstract**—In this paper, recent progress of AlGaIn/GaN-based power high-electron-mobility transistors (HEMT's) is reviewed. Remarkable improvement in performances was obtained through adoption of high Al contents in the AlGaIn layer. The mobility in these modulation-doped structures is about  $1200 \text{ cm}^2 \cdot \text{V}^{-1} \cdot \text{s}^{-1}$  at 300 K with sheet densities of over  $1 \times 10^{13} \text{ cm}^{-2}$ . The current density is over 1 A/mm with gate-drain breakdown voltages up to 280 V.  $f_t$  values up to 52 GHz have been demonstrated. Continuous wave (CW) power densities greater than 3 W/mm at 18 GHz have been achieved.

**Index Terms**—FET, GaN, HEMT, microwave, transistor.

## I. INTRODUCTION

THE NEED for high-power solid-state amplifiers has exploded with the development of wireless communications. Requirements include attributes of ultrahigh power, high efficiency, linearity, manufacturability, and ultimately, low cost. Though the relative importance of each of these features is application specific, it is clear that current technologies fall short of satisfying a majority of the needs simultaneously. Furthermore, market needs are currently satisfied with a multitude of technologies, including silicon- and III-V-based solid-state solutions and the ever-advancing and increasingly reliable vacuum tube solution. Therefore, a new technology has to a large extent a difficult task, that of providing a better solution, not a unique one. Wide-bandgap semiconductors promise the potential of a single technology, which either aids conventional implementations or creates affordable new configurations. In particular, the GaN-based family of semiconductors is attractive, as it has technological and cost advantages over competing semiconductors (e.g., SiC). Firstly, GaN possesses attractive electronic material properties such as a large bandgap (3.4 eV), high breakdown field ( $3 \times 10^6 \text{ V} \cdot \text{cm}^{-1}$ ), the existence of modulation-doped AlGaIn/GaN structures with attendant high electron mobility ( $1500 \text{ cm}^2 \cdot \text{V}^{-1} \cdot \text{s}^{-1}$ ) and extremely high peak ( $3 \times 10^7 \text{ cm} \cdot \text{s}^{-1}$ ) and saturation ( $3 \times 10^7 \text{ cm} \cdot \text{s}^{-1}$ ) [1]. Secondly, it has the technology development costs amortized over several large electronic and opto-electronic applications and, therefore, could potentially be a low-cost solution. In this paper, we will summarize the technology development and status of GaN-based microwave/millimeter-wave power devices.

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## II. EPITAXIAL-MATERIALS DESIGN AND PROPERTIES

Microwave devices in the GaN-based system are presently based on designs developed in the conventional III-V semiconductors, namely the MESFET, heterostructure MESFET, and high-electron-mobility transistor (HEMT) or MODFET. The materials occur in either hexagonal or cubic forms, with the former being stable and the latter meta-stable. Hexagonal crystals are grown on sapphire or 6H/4H SiC, whereas the cubic forms are grown on polar cubic materials such as GaAs. The work on microwave transistors has currently been confined to hexagonal materials. The epitaxial materials are grown either by molecular-beam epitaxy (MBE) or metal-organic-chemical-vapor deposition (MOCVD). Of the possible device structures, the GaN MESFET, though simple to grow and process, does not afford the full advantage of the material, as the effective Schottky-barrier height to doped GaN is relatively low ( $<1 \text{ V}$ ) and, hence, is not preferred in anticipated high-voltage applications. AlGaIn/GaN structures are attractive as both a high-mobility and high-carrier density, through a two-dimensional electron gas (2-DEG), are generated, and a high Schottky-barrier height to AlGaIn is afforded. It is for this reason that most advances have been based on AlGaIn/GaN devices. Recent intensive research on AlGaIn/GaN HEMT's has resulted in a steady increase in power density from 1.1 W/mm at 2 GHz [2] to 1.5–1.57 W/mm at 4 GHz [3], [4] and 1.7 W/mm at 10 GHz [5]. These devices utilized AlGaIn layers with Al mole fractions of 15%–17.5%, as the best 2-DEG mobility is generally achieved with an Al mole fraction of approximately 15%. Little had been done to investigate the impact of a higher Al mole fraction on power performance.

A higher Al mole fraction results in a higher bandgap of the AlGaIn and, hence, a much higher composite breakdown field than the already wide-bandgap GaN. Also, the resultant larger conduction-band discontinuity ( $\Delta E_c$ ) improves carrier confinement, allowing a high mobility to coexist with a large carrier density. In addition to the conventional increase in  $\Delta E_c$  in the AlGaIn system with increasing Al content, an additional benefit exists in the hexagonal-based materials in that the positive piezo-electric charge induced in the AlGaIn at the interface enhances the confining field, allowing increased charge confinement. The saturation velocity  $v_s$ , related to the carriers in the GaN channel, remains high, with little dependence on the AlGaIn layer. Finally, as mentioned earlier, a higher Schottky gate barrier resulting from a higher Al content will more effectively suppress thermionic gate leakage current at elevated operation temperatures. These arguments predict higher equivalent figures of merit with a higher Al content. This has been experimentally demonstrated through improved dc and RF performances.

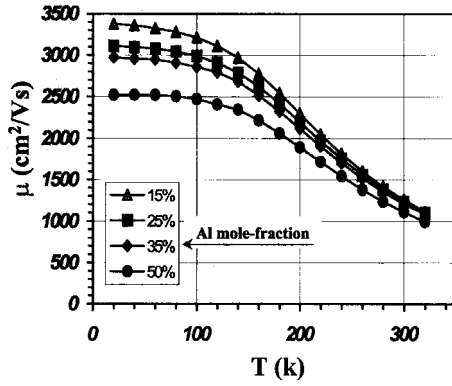


Fig. 1. 2-D gas mobility as a function temperature for various Al compositions.

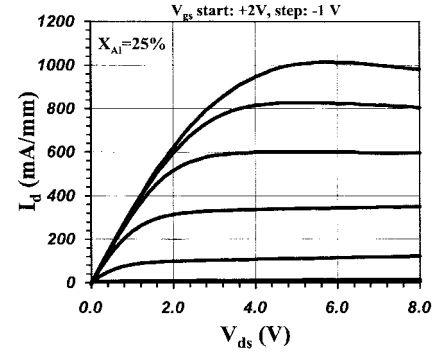
The epi-films under study were grown on *C*-plane sapphire substrates by MOCVD and had the same layer structure: a 200-Å GaN nucleation layer, 2-μm insulating GaN (i-GaN) and a 200-Å AlGa<sub>0.5</sub>N donor-barrier layer. The AlGa<sub>0.5</sub>N layer consisted of a 30-Å unintentionally doped (UID) spacer, a 150-Å Si-doped region, and a 20-Å UID cap. The Al mole fractions of the AlGa<sub>0.5</sub>N were chosen as 15%, 25%, 35%, and 50%. Theoretically, a higher Al content leads to a larger  $\Delta E_c$  and allows a higher 2-DEG density. Experimentally, the highest sheet densities achieved were  $8 \times 10^{12} \text{ cm}^{-2}$ ,  $1 \times 10^{13} \text{ cm}^{-2}$ ,  $1.2 \times 10^{13} \text{ cm}^{-2}$ , and  $1.2 \times 10^{13} \text{ cm}^{-2}$ , respectively. The inability to increase carrier density with  $X_{\text{Al}}$  greater than 35% is attributed to the reduced doping efficiency by Si.

Mobilities as a function of temperature are shown in Fig. 1. Each measured-sheet carrier density was essentially constant through out the temperature range of 20–320 K. The low-temperature (20 K) mobility does reduce with increasing Al content, which may be due to either the interface-roughness scattering enhanced by the higher piezo electric charge density, or the enhanced remote alloy scattering. The 300-K mobilities, however, are nearly the same, indicating that phonon scattering is dominant. For devices operating at room temperature and above, a high Al content up to 50% should introduce little increase in channel resistance.

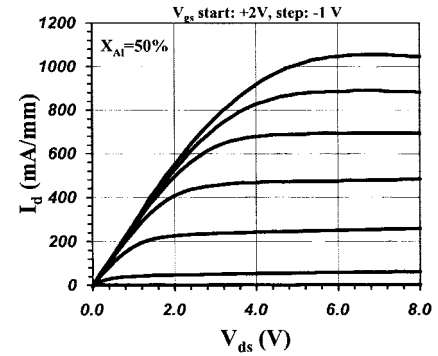
### III. DEVICE PERFORMANCE

Devices were fabricated on epi-films with Al contents of 25% and 50%. The fabrication process was a conventional mesa isolated process with chlorine reactive-ion etching (RIE) used for mesa definition. Ti/Al/Ni/Au was used for ohmic contact formation and Ni/Au for the gate metal. Transfer ohmic contact resistances were measured as 0.5 and 1 Ω · mm, respectively. The higher contact resistance for the latter Al<sub>0.5</sub>Ga<sub>0.5</sub>N/GaN HEMT is attributed to the high Al composition of the layer. The corresponding gate lengths of the devices were 0.9 and 0.7 μm, both obtained by optical lithography. Fig. 2 shows the output *I*–*V* characteristics of both devices.

With the higher charge density of  $1 \times 10^{13} \text{ cm}^{-2}$  compared to the  $5 \sim 7 \times 10^{12} \text{ cm}^{-2}$  value for the Al<sub>0.15</sub>Ga<sub>0.85</sub>N/GaN devices [4], the Al<sub>0.25</sub>Ga<sub>0.75</sub>N/GaN HEMT exhibited a largely increased current density of 1000 mA/mm from the previous



(a)



(b)

Fig. 2. DC *I*–*V* characteristics of AlGa<sub>0.5</sub>N/GaN HEMT's with 25% and 50% Al mole fractions.

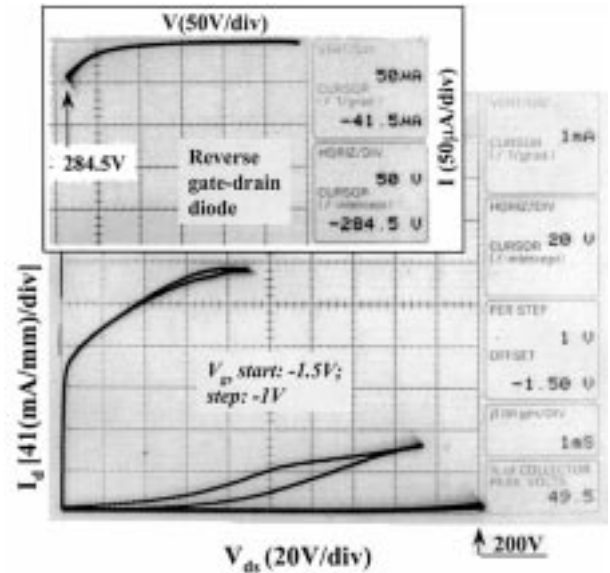


Fig. 3. High-voltage characteristics of an Al<sub>0.5</sub>Ga<sub>0.5</sub>N/GaN HEMT (gate length = 0.7 μm, gate-drain spacing = 3 μm).

peak of 700 mA/mm. The on resistance was also reduced to 3 Ω · mm, resulting in a knee voltage of only 5 V for such a high current density. The much larger peak transconductance of 255 mS/mm than the previous 160 mS/mm is attributed to both the thinner AlGa<sub>0.5</sub>N layer and the reduced access resistances. The Al<sub>0.5</sub>Ga<sub>0.5</sub>N/GaN HEMT mildly suffered from the poor

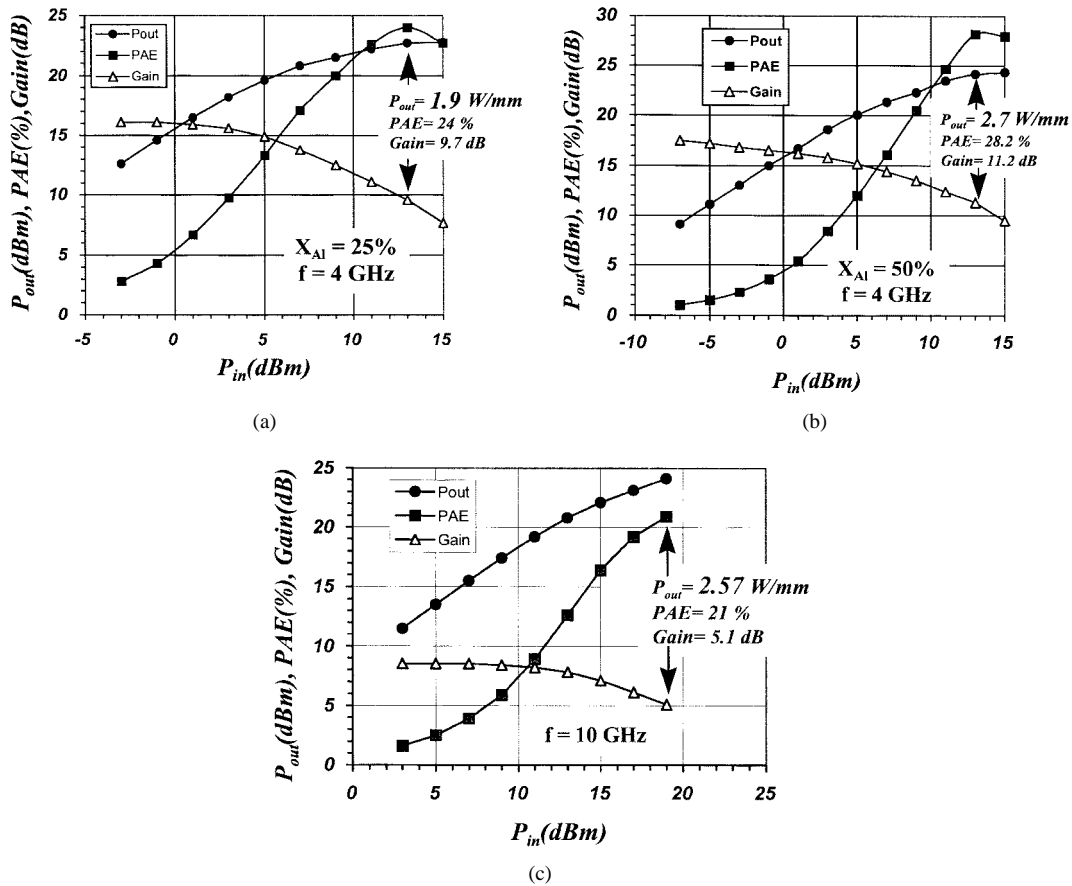


Fig. 4. Microwave power performance of the  $\text{Al}_x\text{Ga}_{1-x}\text{N}/\text{GaN}$  HEMT's by optical lithography. (a)  $X_{\text{Al}} = 25\%$ ,  $f = 4$  GHz. (b)  $X_{\text{Al}} = 50\%$ ,  $f = 4$  GHz. (c)  $X_{\text{Al}} = 50\%$ ,  $f = 10$  GHz. Device width: 100  $\mu\text{m}$ .

ohmic contact resistance. Nevertheless, with a charge density of  $1.2 \times 10^{13} \text{ cm}^{-2}$ , it demonstrated a current density of 1050 mA/mm. The knee voltage is 6 V and transconductance 220 mS/mm.

Despite the much higher charge density over the previous  $\text{Al}_{0.15}\text{Ga}_{0.85}\text{N}/\text{GaN}$  devices [4], the  $\text{Al}_{0.25}\text{Ga}_{0.75}\text{N}/\text{GaN}$  HEMT's showed similar gate-drain breakdown voltages around 100, 160, and 220 V for gate-drain separations of 1, 2, and 3  $\mu\text{m}$ , respectively. The corresponding values for the  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{N}/\text{GaN}$  devices were generally 20% higher. Fig. 3 is the high-voltage  $I$ - $V$  characteristics of an  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{N}/\text{GaN}$  HEMT with a gate-drain spacing of 3  $\mu\text{m}$ , showing a gate-drain breakdown  $>284$  V and a three-terminal voltage  $>200$  V before punching through. These voltages and the current density of 1 A/mm correspond to an ultrahigh  $I$ - $V$  product per unit width of  $(I_{\text{max}}V_{\text{max}}) > 200 \text{ VA/mm}$ . While a maximum power density of  $(I_{\text{max}}V_{\text{max}})/8 = 25 \text{ W/mm}$  is calculated, further work needs to be done to verify this potential.

Small-signal RF measurements yielded current-gain and power-gain cutoff frequencies ( $f_t$  and  $f_{\text{max}}$ ) of 15 and 35 GHz for the 0.9- $\mu\text{m}$  gate-length  $\text{Al}_{0.25}\text{Ga}_{0.75}\text{N}/\text{GaN}$  devices, 17.5 and 44 GHz for the 0.7- $\mu\text{m}$  gate-length  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{N}/\text{GaN}$  devices, respectively. Corresponding intrinsic  $f_t$ 's of 17.2 and 21.8 GHz were extracted from the  $S$ -parameters, translating to intrinsic  $f_t$  gate-length products of 15.5 and 15.3-GHz  $\mu\text{m}$ , considerably improved from the previous value of 11-GHz  $\mu\text{m}$  [4].

Continuous wave (CW) microwave power performance was characterized un-cooled on the sapphire substrate with a Maury load-pull system. As seen in Fig. 4(a), at 4 GHz, the  $\text{Al}_{0.25}\text{Ga}_{0.75}\text{N}/\text{GaN}$  HEMT's exhibited moderately improved output power density and power-added efficiency (PAE) of 1.9 W/mm and 24%, respectively, from the previous values of 1.57 W/mm and 20% [4]. In contrast, the  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{N}/\text{GaN}$  devices produced a largely enhanced power density of 2.7 W/mm with a PAE of 28% and a large-signal gain of 11.2 dB [Fig. 4 (b)]. Fig. 4(c) is the measurement result of an  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{N}/\text{GaN}$  HEMT at 10 GHz, showing a power density of 2.57 W/mm with a PAE and a large-signal gain of 21% and 5.1 dB, respectively. At 8 GHz, the corresponding values were 2.84 W/mm, 23% and 6.6 dB.

The CW power density of 2.57–2.84 W/mm up to X-band is a remarkable improvement over previous GaN-based FET's in literature. An important observation is that the current of the  $\text{Al}_{0.25}\text{Ga}_{0.75}\text{N}/\text{GaN}$  HEMT's degraded much more than the  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{N}/\text{GaN}$  devices after the high-power stress up to 8.5 W/mm, which may be related to the interaction between the gate-metal and AlGaN layer. The result will be discussed in a later publication. Fig. 5 summarizes the best power results measured un-cooled on sapphire substrates for various Al mole fractions. The power density monotonically increases with  $X_{\text{Al}}$ , supporting the benefit of using higher Al contents.

To take the advantage of a higher electron velocity, short gate-length  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{N}/\text{GaN}$  devices were also fabricated.

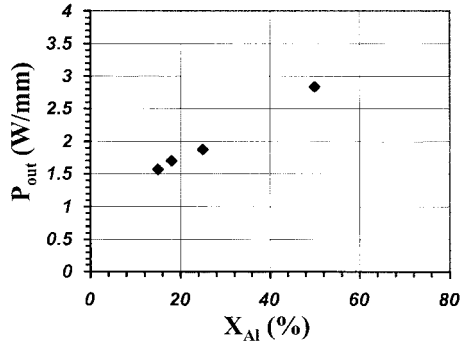


Fig. 5. Best measured output-power density versus Al mole fraction for AlGaIn/GaN HEMT's by optical lithography (uncooled on sapphire substrates).

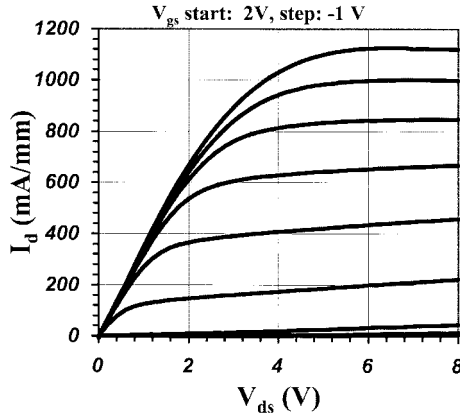


Fig. 6.  $I$ - $V$  characteristics of a 0.25- $\mu\text{m}$  gate-length  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{N}/\text{GaN}$  HEMT.

The gate definition was by electron-beam lithography with a routine T-gate technology. The gate length was 0.25–0.28  $\mu\text{m}$ , while the gate–source and gate–drain separations were 0.6 and 0.75  $\mu\text{m}$ , respectively. Ohmic contact resistance was measured as 0.9–1  $\Omega \cdot \text{mm}$ , confirming previous result that the contact resistance of the  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{N}/\text{GaN}$  HEMT's is higher than the 0.5  $\Omega \cdot \text{mm}$  value generally achieved with lower Al-content devices. Fig. 6 is the output  $I$ - $V$  characteristics showing a high-saturation current density of 1130 mA/mm and a maximum extrinsic transconductance of 240 mS/mm. Both specifications are better than the 0.7- $\mu\text{m}$  gate-length devices due to the shorter channel (which results in a channel velocity closer to saturation). The on-resistance is 2.85  $\Omega \cdot \text{mm}$  where the source resistance accounts for about 1.3  $\Omega \cdot \text{mm}$ . The gate–drain breakdown voltage is around 80 V. Small-signal microwave characterization performed on devices with 100- $\mu\text{m}$  gatewidth yielded  $f_t$  and  $f_{\text{max}}$  of 52 and 82 GHz, respectively, as shown in Fig. 7. The  $f_t$  represents the highest reported to date for a GaN-channel FET, while the relatively low  $f_{\text{max}}/f_t$  ratio is related to the high ohmic contact resistance. An active load–pull system was used for CW power measurements. The output tuning was performed on the fundamental frequency only. The input of the device was not tuned and was connected to the microwave source with a 50- $\Omega$  line. The actual input power was calculated with the available input power and the measured input reflection coefficient

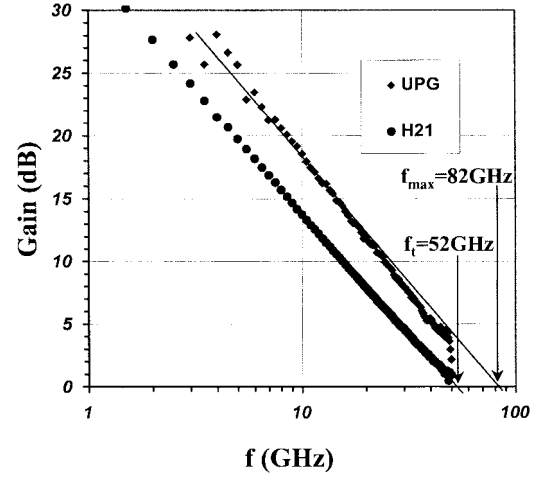


Fig. 7. Current and power gain versus frequency for a 0.25- $\mu\text{m}$  gate-length  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{N}/\text{GaN}$  HEMT (device width: 100  $\mu\text{m}$ , bias: drain current = 300 mA/mm, and a source–drain voltage = 8 V).

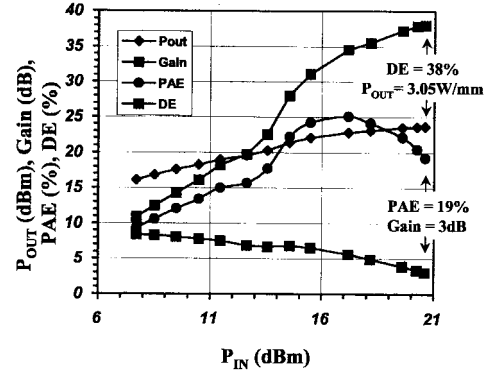


Fig. 8. Power performance of a 0.25- $\mu\text{m}$  gate-length  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{N}/\text{GaN}$  HEMT at 18 GHz (device dimension: 0.25  $\mu\text{m} \times 76 \mu\text{m}$ ).

( $S_{11}$ ). Fig. 8 shows the measurement result at 18 GHz for a device with a gatewidth of 76  $\mu\text{m}$ , biased in class-AB mode with a drain voltage of 21 V and a quiescent drain current of 190 mA/mm (which was self-adjusted to 380 mA/mm at the maximum input drive). The small-signal linear gain is  $\sim 7.5$  dB. The peak output power is 23.65 dBm, which is normalized to a power density of 3.05 W/mm. The corresponding large-signal gain, PAE, and drain efficiency (DE) are 3.1 dB, 19.2%, and 38%, respectively. Many devices with gatewidth of 76 and 100  $\mu\text{m}$  showed output power densities above 3 W/mm, while the highest measured was 3.3 W/mm with a large-signal gain, PAE, and DE of 2.4 dB, 18.2%, and 42.7%, respectively. Such CW power densities are three times as high as generally achieved with GaAs MESFET's and are two times as high as those for the AlGaIn/GaN devices with lower Al-contents of 15%–17% [3]–[5]. Compared with recent 4H SiC MESFET's with similar power densities at 10 GHz under pulsed condition [6], the  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{N}/\text{GaN}$  HEMT's showed a higher operation frequency.

The effect of device width on power performance was also investigated at 8 GHz, as shown in Fig. 9. It is seen that although output power increases with increasing gatewidth,

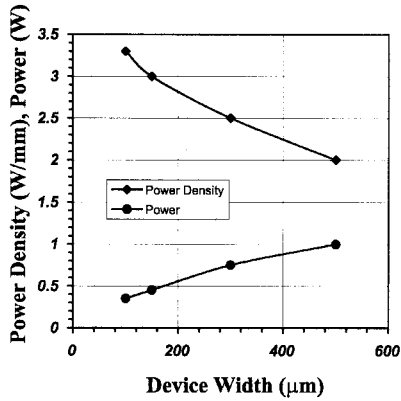


Fig. 9. Power density and total output power versus gatewidth for the 0.25- $\mu\text{m}$  gate-length  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{N}/\text{GaN}$  HEMT's at 8 GHz.

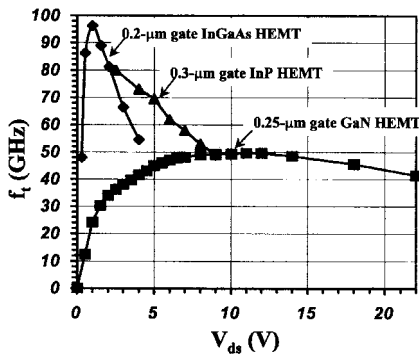


Fig. 10. Current gain cutoff frequencies versus drain bias for an  $\text{Al}_{0.5}\text{Ga}_{0.5}\text{N}/\text{GaN}$  HEMT (gate bias:  $V_{gs} = -3$  V) as compared with GaInAs and InP channel HEMT's with similar gate lengths [7], [8].

the power density actually reduces. This is attributed to more severe self-heating for larger devices. Nevertheless, a total output power of 1 W is obtained at X-band with a 500- $\mu\text{m}$ -wide device.

Finally, as a reason explaining the superior power capability at high frequencies for the GaN-channel devices, a compelling comparison, as shown in Fig. 10, is presented between the drain-bias dependence of AlGaIn/GaN HEMT's versus conventional InGaAs- and InP-based HEMT's [7], [8]. As power applications require that the bias be high, it is necessary that the RF figures of merit remain high at large biases. It is clear that the  $f_T$ 's of the conventional devices collapse at increased drain bias because of the drain delay in the extended depletion region, whereas  $f_T$  is maintained high in the GaN system. Though not fully understood, this suggests that the high saturation velocity of electrons in GaN may be mitigating the drain delay.

#### IV. CONCLUSIONS

Remarkable progress of both the AlGaIn/GaN materials quality and device performance has been achieved. Within less than two years, the power density has been improved from 1.1 W/mm at 2 GHz [2] to over 3 W/mm at 18 GHz [9]. A total power of 1 W at X-band has also been demonstrated. However, several fundamental issues are unresolved. What is the nature of the GaN and AlGaIn surface, why is the

breakdown field strength in the highly doped two-dimensional HEMT structure approaching that predicted in undoped one-dimensional GaN, what is the electron transport physics in the HEMT and what is the best manner in which to thermally manage the device, are just a few of the questions. Nevertheless, a material system that demonstrates the high-power capability of  $>3$  W/mm at 18 GHz in a comparatively immature stage deserves the attention to answer the questions.

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